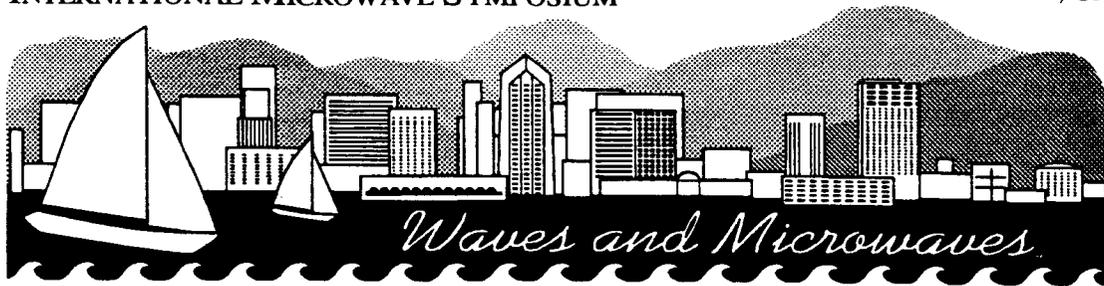


1994 IEEE MTT-S  
INTERNATIONAL MICROWAVE SYMPOSIUM

MAY 23-27, 1994  
SAN DIEGO, CA



# WORKSHOP NOTES

## WMHE

**Computer-Aided Design of  
Superconducting Microwave  
Components.**

Monday, May 23, 1994

8:00AM - 12:00PM

San Diego Convention Center, Room 9





**Approximate Boundary Conditions  
For Modeling Superconducting Thin Film  
Microwave Circuits**

presented at:

**CAD OF SUPERCONDUCTING MICROWAVE COMPONENTS  
WORKSHOP WHHE**

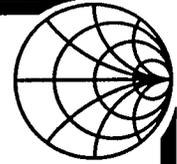
**1994 IEEE International Microwave Symposium  
San Diego, CA**

**Monday, 23 May 1994**

by:

**Jeffrey M. Pond, Code 6851  
Microwave Technology Branch  
Naval Research Laboratory  
Washington, DC 20375**

65



## What Are Approximate Boundary Conditions And Why Use Them?

Absolute Boundary Conditions are:

An absolute value for a component of the electromagnetic field is specified such as:

Neumann condition

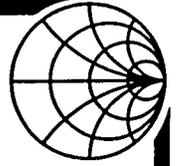
Dirichlet condition

Approximate Boundary Conditions are:

Approximations to field matching conditions which, by invoking an appropriate limit, reduce a  $n$ -region problem to an  $(n-1)$ -region problem thereby reducing computational difficulty. Examples are:

Surface Impedance (Leontovich) - reduces 2 region to 1 region

Resistive Sheet - reduces 3 region to 2 region



## What are the approximations?

### Surface Impedance Boundary Condition:

Propagating wave in second region must be attenuated to zero, i.e. semi-infinite or many skin/penetration depths thick

Angle of refraction must be effectively zero for any angle of incidence (implies a 'dense' medium)

Boundary condition given by ratio of tangential electric field to tangential magnetic field as an impedance:  $E_{\text{tan}} = Z_S J_S = Z_S \mathbf{n} \times \mathbf{H}$

### Resistive Sheet Boundary Condition:

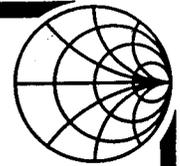
Eliminate electrically thin region with the field matching conditions:

$$\mathbf{n} \times (\mathbf{E}^+ - \mathbf{E}^-) = 0$$

$$\mathbf{n} \times (\mathbf{H}^+ - \mathbf{H}^-) = \mathbf{J}_S = -(1/R) \mathbf{n} \times (\mathbf{n} \times \mathbf{E}^{\pm})$$

$$R = 1/(\sigma t)$$

67



## Application to Superconductors

Surface Impedance Boundary Condition:

Accurate if the superconductor is much thicker than a penetration depth

Resistive Sheet Boundary Condition:

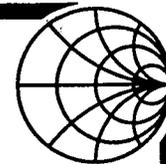
Accurate if the superconductor is much thinner than a penetration depth

$$\sigma = \sigma_n - j \sigma_s = \sigma_n - j (\omega \mu_0 \lambda^2)^{-1}$$

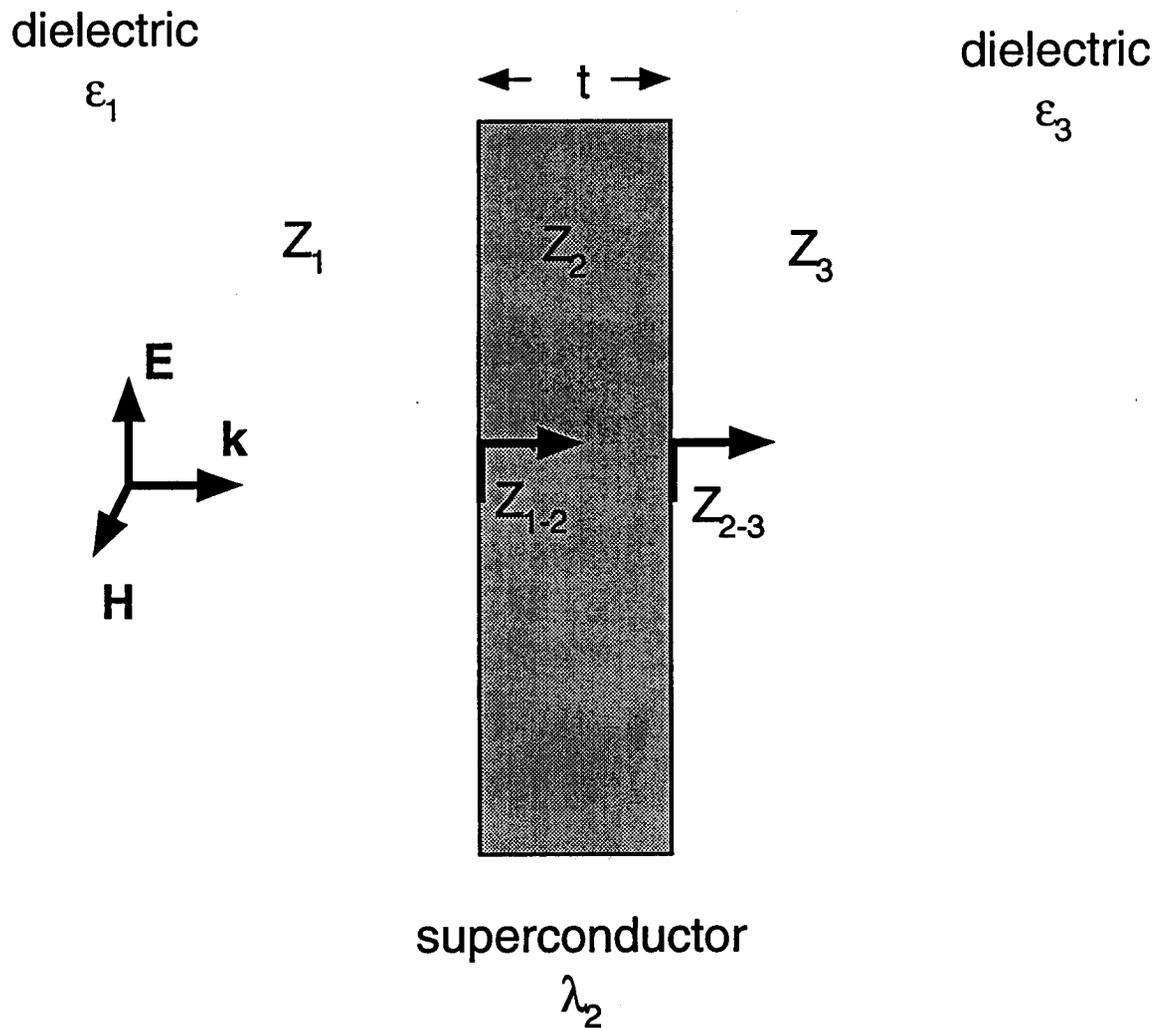
Generalized Resistive Sheet Boundary Condition:

1st order correction for films where the penetration depth is on the order of the sheet thickness

$$R = Z_s \coth (t/\lambda)$$

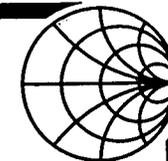


# Effective Surface Impedance of a Finite Thickness Superconductor



69





## Effective Surface Impedance of a Finite Thickness Superconductor (cont.)

Under what conditions does:

$$Z_{\text{eff}} = Z_{1-2} = Z_2 \frac{Z_2 + Z_{2-3} \coth(k_2 t)}{Z_{2-3} + Z_2 \coth(k_2 t)}$$

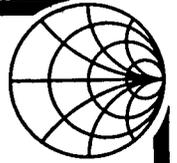
reduce to:

$$Z_{\text{eff}} = Z_{1-2} \cong Z_2 \coth(k_2 t) \quad ?$$

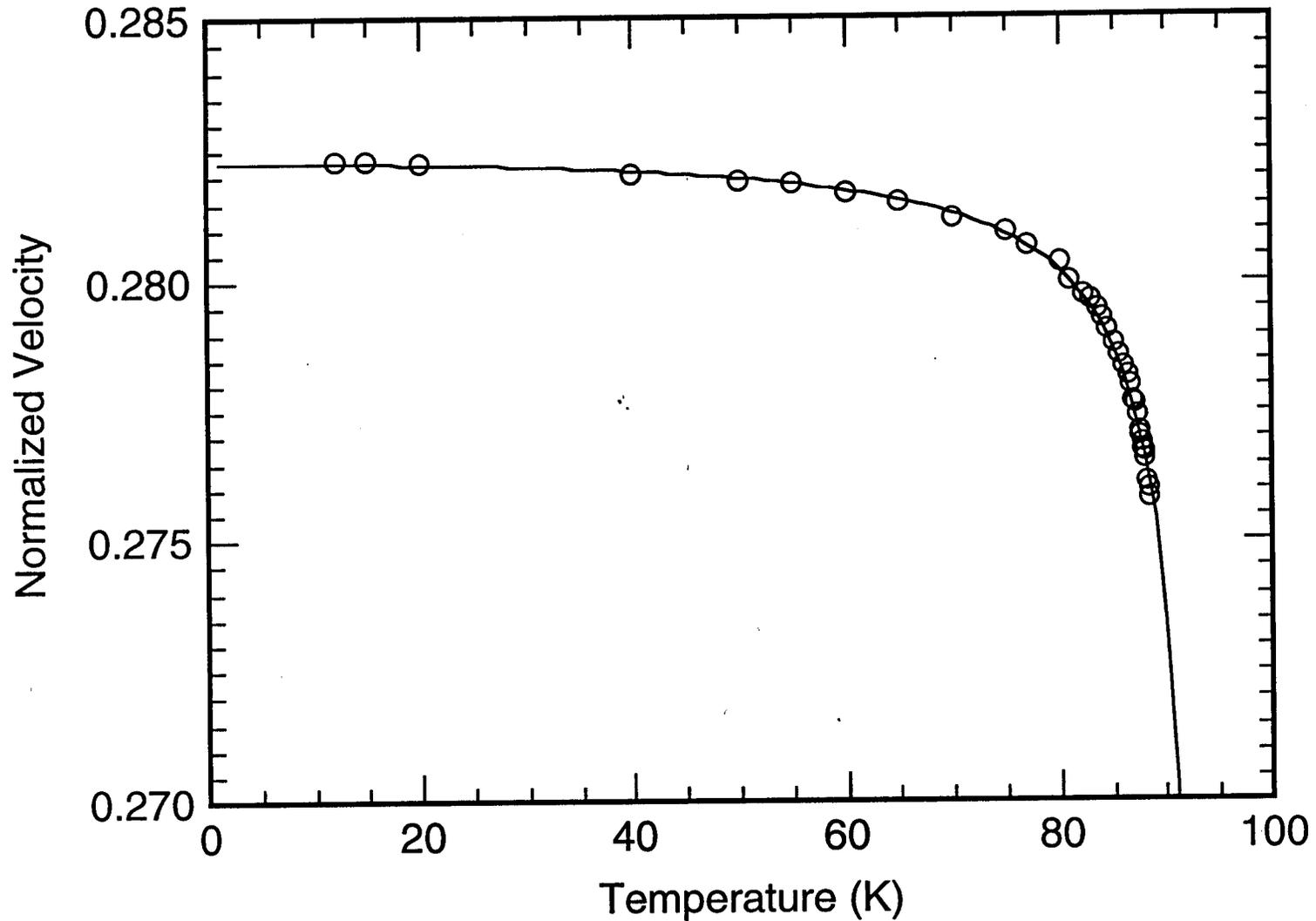
When:

$$|Z_{2-3} \coth(k_2 t)| \gg |Z_2| \quad \text{and} \quad |Z_{2-3}| \gg |Z_2 \coth(k_2 t)|$$

which is easily satisfied if  $Z_{2-3} = Z_3$  (a semi-infinite dielectric region) for all practical superconductor thicknesses.



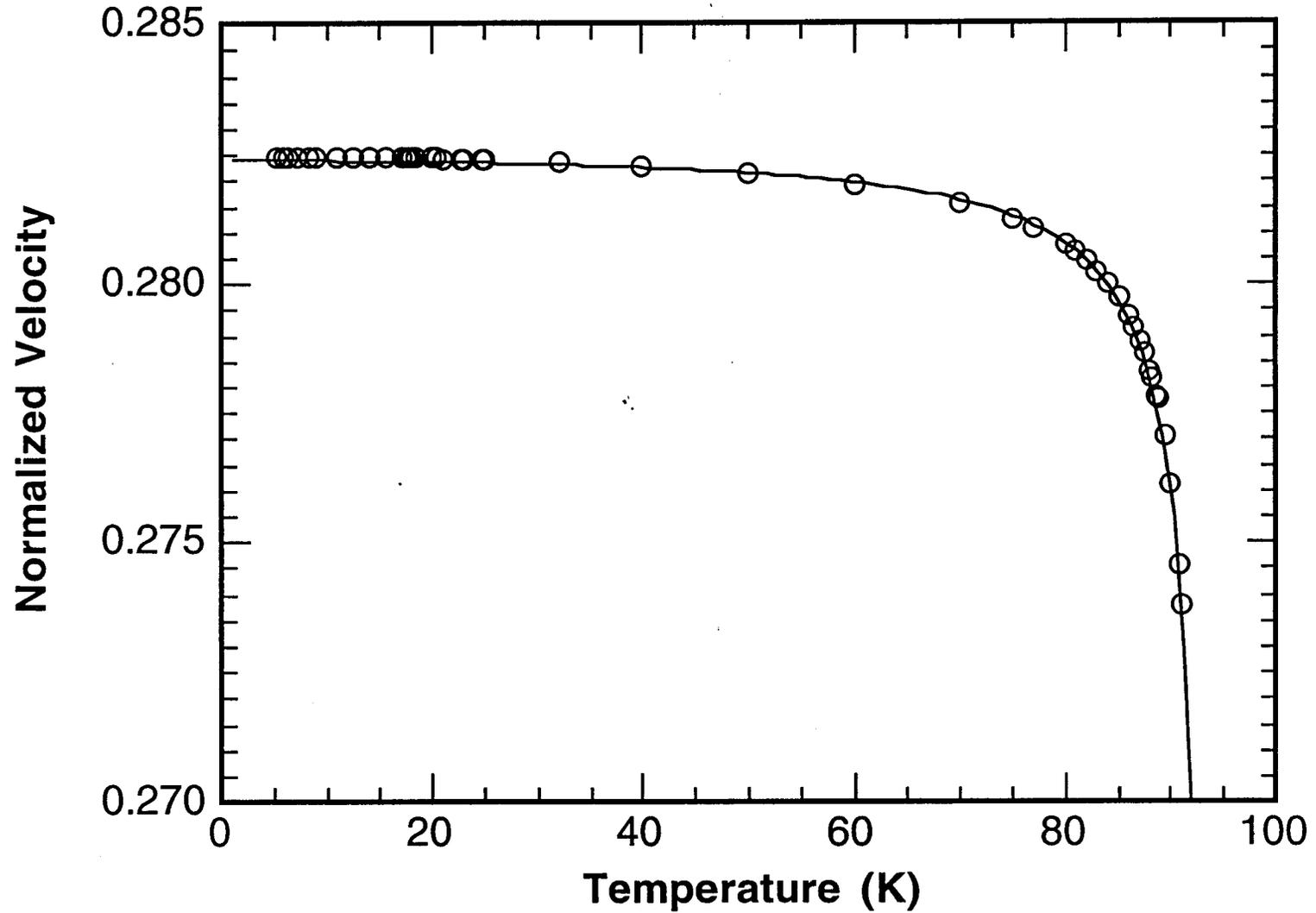
**Measurement and Spectral-Domain/Approximate-Boundary-Condition Model for a Coplanar Waveguide Fabricated from 100-nm thick YBCO**



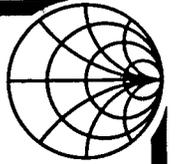
71



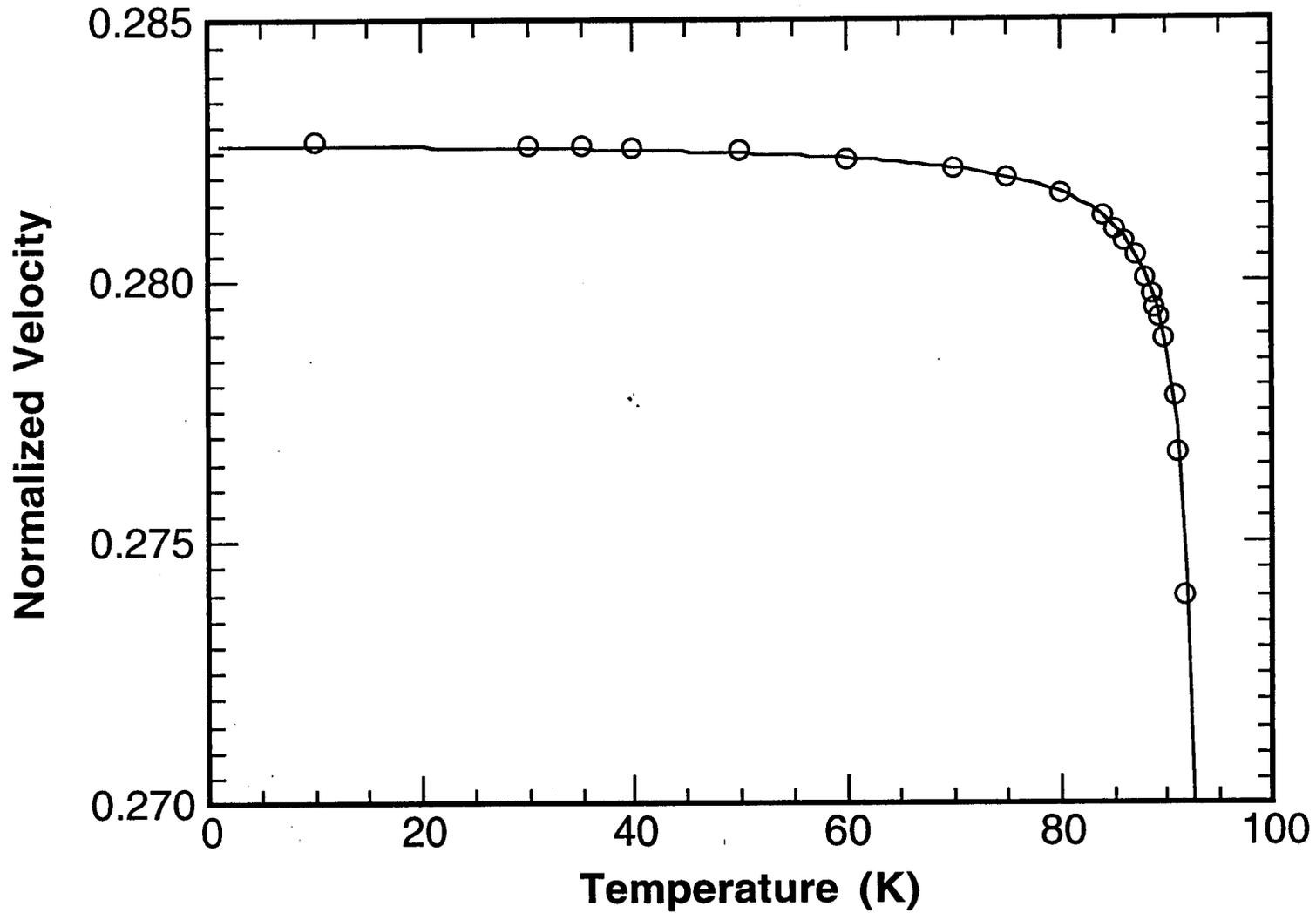
### Measurement and Spectral-Domain/Approximate-Boundary-Condition Model for a Coplanar Waveguide Fabricated from 200-nm thick YBCO



72



**Measurement and Spectral-Domain/Approximate-Boundary-Condition Model for a Coplanar Waveguide Fabricated from 300-nm thick YBCO**



72



## Low Temperature Self-consistency for a Given Geometry

Phase velocity is given by:

$$v_p = \frac{1}{\sqrt{L_e C \left(1 + \frac{L_i}{L_e}\right)}} = \frac{1}{\sqrt{L_e C \left(1 + Gt \left(\frac{\lambda}{t}\right) \coth\left(\frac{t}{\lambda}\right)\right)}}$$

At low temperature the normalized phase velocity difference is given by:

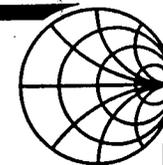
$$-\frac{v_p - v_{p0}}{c} = \frac{v_{p0}}{c} \frac{Gt}{2} \left[ \frac{\lambda_0}{t} \coth\left(\frac{t}{\lambda_0}\right) + \operatorname{csch}^2\left(\frac{t}{\lambda_0}\right) \right] f(T)$$

where:

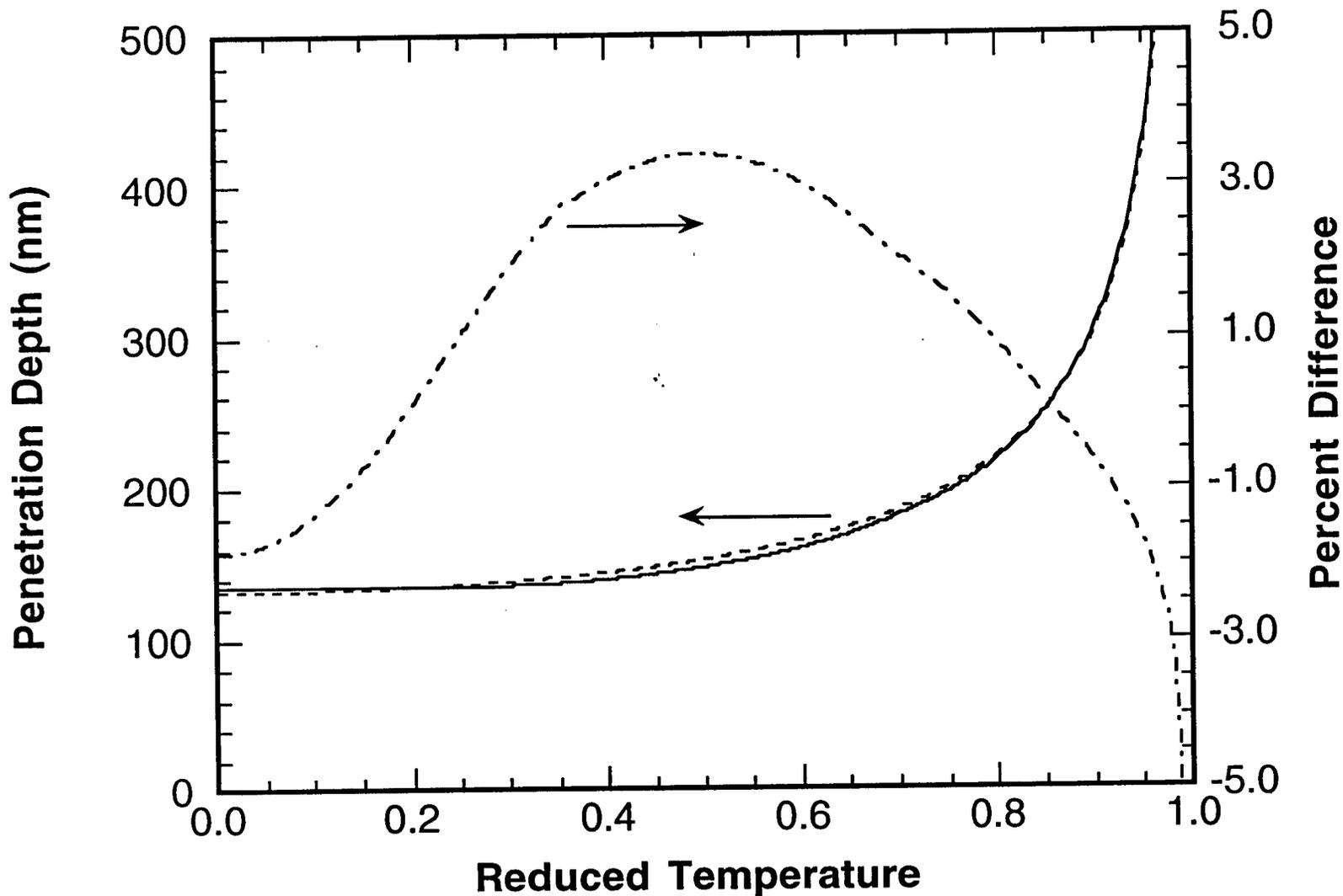
$$\lambda = \frac{\lambda_0}{\sqrt{1 - f(T)}}$$

The value of  $G$  had better be the same for all devices patterned with the same mask. In fact the values obtained for  $G$ , after obtaining a best fit over the entire temperature range, were within  $\pm 3\%$

74



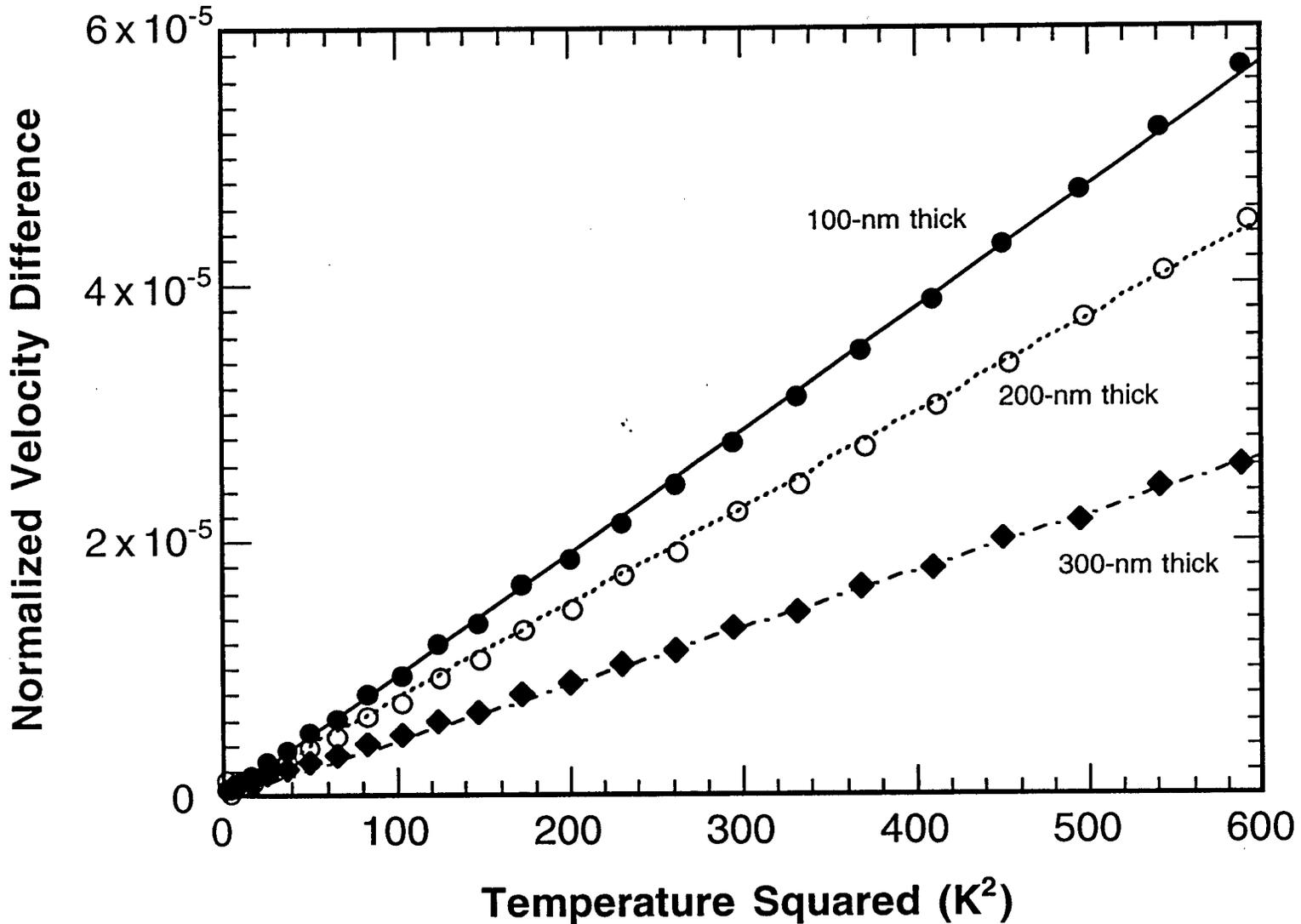
## Comparison of the BCS Temperature Dependence of the Penetration Depth to a $(1-(T/T_c)^2)^{-1/2}$ Dependence



75



# Low Temperature Normalized Phase Velocity Difference of the Three Coplanar Waveguide Transmission Lines



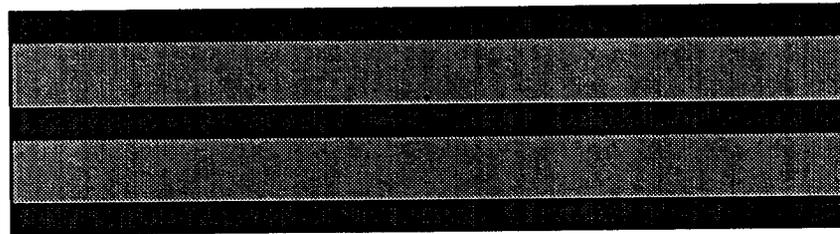
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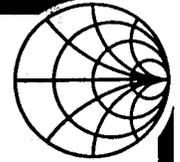
## Effective Surface Impedance of a Finite Thickness Superconductor (cont.)

Under what practical situations does the hyperbolic cotangent correction fail?

Consider a set of inductively-coupled parallel plate waveguides:



Solving for the two propagating modes by treating the center superconductor as a resistive sheet with a hyperbolic cotangent correction will result in an the wrong answer for both modes.



## Limitations of Hyperbolic Cotangent Correction and Possible Improvements

Assumes a particular field distribution inside the superconductor

Can fail when other very low loss conducting surfaces are in close proximity - must be used with care

Need is for a robust approximate boundary condition that is more universally valid for thicknesses on the order of a penetration depth

Modification by the Green's function of the external region may address these issues - investigation in progress

Modification to include higher-order field matching, i.e. describe discontinuities in field derivatives in terms of superconductor parameters